

# Curvature of the horizon

*Abstract:* Followers of the flat Earth idea frequently use quantities like distance to the horizon and its dip. Mathematical expressions for these can be easily derived and formulae for good practical approximations are well known in literature. For example the dip in arc minutes is calculated as  $1.78 \cdot \sqrt{h}$ , where  $h$ , the height of the observer, is expressed in meters. Such formulae are also derived in this article but what is new here is the formula for curvature of the horizon as seen from the perspective of an observer. If the observer sees a part of the horizon contained within the angle  $\gamma$  then the horizon arc in the middle of this field of view will be seen higher than the ends of the arc (i.e. above the chord) by the angle (again in arc minutes):

$$g = 1.78 \cdot \sqrt{h} \cdot [1/\cos(\gamma/2) - 1].$$

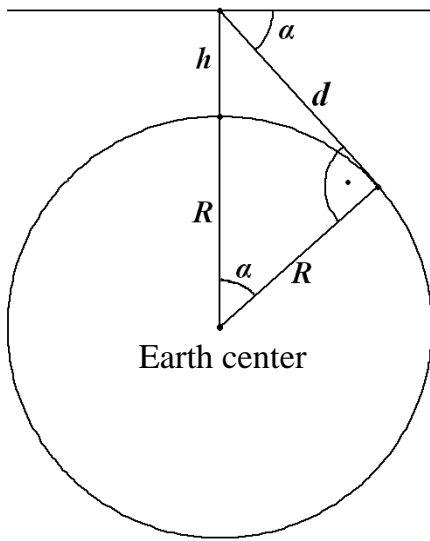
This is just the dip scaled by the factor  $1/\cos(\gamma/2) - 1$ . This formula proves quite accurate and thus it can be safely used by all flat earthers. Assessment of its accuracy can be made by comparing the results it gives with the results obtained from exact solution as given in the Table 2 below. For example, this formula errs by less than 1% for  $h \leq 100\,000$  m and  $\gamma = 90^\circ$ , or  $h \leq 50\,000$  m and  $\gamma = 120^\circ$ , or  $h \leq 10\,000$  m and  $\gamma = 150^\circ$ . Actual observations presented in [this web site](#) are shown to agree with our theoretical calculations.

Flat earth believers often cite the lack of evidence for the earth's sphericity, in particular the fact that the curvature of the horizon is not visible on a daily basis, even at considerable heights. I recently watched a video by someone from Poland who planned to send a balloon with equipment high enough to show that the curvature is not visible there either. Stratospheric balloons usually rise to an altitude of 20-40 km (the record being 43,561 m).

The difficulty in seeing the curvature with the naked eye or in photographs is simply due to the smallness of the curvature. Nevertheless, the fact that the horizon is lowered has been repeatedly verified experimentally. A good example of ground measurements with a theodolite is presented in the [video here](#).

The formulas for calculating the distance to the horizon and its lowering or dip relative to the plane perpendicular to the vertical at the observer's location are quite simple – they can be easily found in the literature. However, calculating the expected amount of curvature of the horizon in photographs of a section of the horizon is somewhat challenging, which prompted me to develop the precise solutions presented here. They will allow for more reliable planning of possible attempts to experimentally confirm or disprove the sphericity of the Earth.

## The distance and dip of the horizon



The distance and dip of the horizon can be easily determined based on the figure on the left. If we assume that the Earth's surface is a sphere with radius  $R$ , the horizon will be a circle with radius  $r = R \cdot \sin \alpha$ , where  $\alpha$  is the horizon dip. This is also the radius of curvature of the horizon. The horizon dip is the angle between the plane perpendicular to the vertical direction at the observer's location and the direction of the horizon. We can calculate it from the relationship  $\sin(\pi/2 - \alpha) = \cos \alpha = R/(R + h)$ , where  $\pi$  is a mathematical constant (3.14159...), which, when expressing an angle in radians, corresponds to a value of  $180^\circ$ . Thus:

$$\alpha = \arccos \frac{R}{R+h} = \arcsin \left\{ \sqrt{1 - \left( \frac{R}{R+h} \right)^2} \right\}$$

We can calculate the distance to the horizon just as easily using Pythagoras' theorem  $d^2 = (R + h)^2 - R^2 = 2h \cdot R + h^2$ , from which

$$d = \sqrt{2h \cdot R + h^2},$$

or from the already calculated angle  $\alpha$

$$d = R \cdot \tan \alpha,$$

where  $R$ ,  $h$  and  $d$  are expressed in the same units, and  $\alpha$  – in radians.

These formulas are accurate and remain valid even for satellites, but since in practice  $h$  is much smaller than  $R$ , thus  $\alpha$  is usually small, we can use the following approximations for this quantity:

$$\alpha \approx \sin \alpha = d/(R + h) = (2hR + h^2)^{1/2}/(R + h) \approx (2h/R)^{1/2} \equiv \sqrt{2h/R}.$$

Assuming  $R = 6371000$  m (the average radius of our planet) for the Earth and converting the units from radians to arc minutes (factor  $180 \cdot 60/\pi$ ), we get

$$\alpha \approx (180 \cdot 60/\pi) \cdot \sqrt{2/R} \cdot \sqrt{h} = 1.926 \cdot \sqrt{h},$$

where  $h$  is expressed in meters and  $\alpha$  in minutes. This approximation can be found in the literature, but with a factor of 1.8 in front of the square root, which significantly compensates for atmospheric refraction (see below).

Table 1 shows the calculations of these two quantities for several heights  $h$  above the Earth's surface. It shows that the approximation of  $1.926\sqrt{h}$  is very

good up to an altitude of 100 km, where the error is  $100 \cdot (609.1 - 605.2) / 605.2 = 0.64\%$ . Even at an altitude of 1000 km, the error is only about 6%.

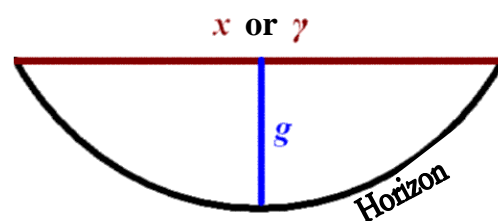
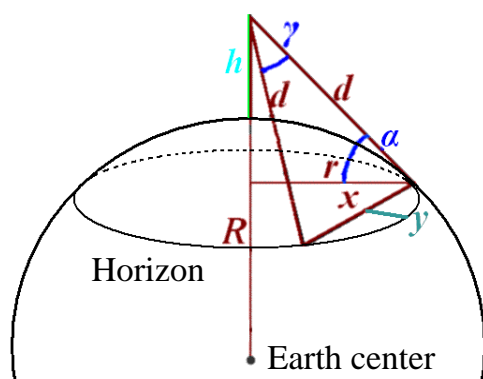
**Table 1.** Horizon dip,  $\alpha$ , and distance,  $d$ , as a function of the observer's height above the Earth,  $h$  ( $r$  is the radius of the horizon), calculated with approximate formula and exactly.

$h$	$1.926 \cdot \sqrt{h}$	$\arccos[R/(R+h)]$	$d$	$r$
[m]	[']	[']	[m]	[m]
10	6.1	6.1	11288	11288
100	19.3	19.3	35696	35695
1000	60.9	60.9	112884	112867
10000	192.6	192.5	357099	356539
25000	304.5	304.1	564955	562747
50000	430.7	429.3	799749	793522
100000	609.1	605.2	1133225	1115713
500000	1361.9	1319.6	2573130	2385884
1000000	1926.0	1811.6	3707020	3204100
10000000	6090.5	4025.9	15080450	5868765
100000000	19260.0	5194.0	106180035	6359562

## Horizon Curvature

When we are positioned exactly at the center of a circle, we see each segment of the circle as its chord, i.e., as a straight line segment. However, when viewing such a segment of circle from any height, we see it as an arc above the chord. The bulge above the chord becomes more pronounced the higher we look from and the larger the angle of view. If we take photos of the horizon from a height  $h$  (see the figures below) and our camera or camcorder has a field of view (angle)  $\gamma$ , the ends of the part of the horizon visible in the resulting image (photo) will be as far from the device (camera) as any point on the entire horizon, i.e., by  $d$ , and the distance between them, i.e., the chord length,  $x$ , will be:

$$x = 2d \cdot \cos(\pi/2 - \gamma/2) = 2(2h \cdot R + h^2)^{1/2} \sin(\gamma/2).$$



The center of the chord connecting the mentioned ends will be at a distance from the observer  $d \cdot \cos(\gamma/2)$ . The distance,  $r'$ , of this center from the vertical line and the depression of the direction towards this center from the camera relative to the horizon plane,  $\alpha'$ , will be respectively:

$$r' = d \cdot [\cos^2(\alpha) - \sin^2(\gamma/2)]^{1/2} \text{ and}$$

$$\alpha' = \arccos[r'/(d \cdot \cos(\gamma/2))] = \arcsin[\sin(\alpha)/\cos(\gamma/2)].$$

The non-obvious right-hand expression for  $\alpha'$  is easily obtained using the relationship  $\arccos x = \arcsin[\sqrt{1 - x^2}]$ .

On the horizon plane, the distance between the cord center and the top of horizon above it is

$$y = r - r' = R \cdot \sin \alpha - r' = d \cdot \cos \alpha - r'.$$

From the vantage point, the top of the horizon bulge will be visible above the chord at an angular distance  $g = \alpha' - \alpha$ . Thus we have an exact expression for this quantity:

$$g = \arcsin[\sin(\alpha)/\cos(\gamma/2)] - \alpha,$$

but for small angles  $\alpha$  (which in practice is true for angles  $\gamma$  less than  $120^\circ$ , even up to a height of  $h = 50,000$  m), we can simplify it to

$$g \approx \alpha/\cos(\gamma/2) - \alpha = \alpha [1/\cos(\gamma/2) - 1]$$

to obtain a simple formula (giving the result in radians):

$$g \approx \sqrt{(2h/R)} \cdot [1/\cos(\gamma/2) - 1].$$

Expressing  $h$  in meters and scaling the whole thing as before for  $\alpha$ , we obtain a novel formula (giving the result  $g$  in minutes of arc):

$$g \approx 1.926 \cdot \sqrt{h} \cdot [1/\cos(\gamma/2) - 1].$$

As a measure of the bulge of the horizon, we can use the percentage ratio of the angle  $g$  to the field of view or the angular distance between the edges of the horizon visible in the image, i.e.,  $100 \cdot g/\gamma$ .

## The Effect of Refraction and Practical Formulas

Refraction (change of direction of light at the boundary of atmospheric layers with different refractive indices, which depend primarily on air temperature and pressure) allows one to see slightly beyond the geometric horizon from altitudes greater than zero – the further beyond the higher is the altitude. At altitudes reaching the stratosphere, a correction for refraction of approximately  $0.5^\circ$  can be expected. In the atmosphere, a light ray follows an upward convex curve. This means that a ray tangent to the spherical Earth at a geometric distance  $d$  from the observer will arrive slightly below the observer (located at altitude  $h$ ), which is equivalent to a reduction in angle  $\alpha$ . In calculations of the horizon distance and dip, this can be taken into account by increasing the Earth's radius,  $R$ . One model compensates for refraction by increasing the Earth's radius by a factor of  $7/6$ , leading to the approximations:

$$\alpha \approx 1.926 \cdot \sqrt{(6/7)} \cdot \sqrt{h} = 1.78 \cdot \sqrt{h}$$

$$d \approx \sqrt{2h \cdot (7/6) \cdot R + h^2} = \sqrt{[(7/3) \cdot h \cdot R + h^2]} \approx 3856 \cdot \sqrt{h}.$$

If we express  $h$  in meters in these formulas, we get  $\alpha$  in arc minutes, and  $d$  in meters. Expressing  $d$  in nautical miles ( $1 \text{ NM} = 1852 \text{ m}$ ;  $3856/1852 = 2.08$ ), we obtain the formula used in navigation:  $d = 2.08 \cdot \sqrt{h}$  (e.g., [Navipedia](#)). However, these approximations fail at high altitudes, say above 100 km. Refraction can be similarly accounted for in the formula for the bulge  $g$ :

$$g \approx \sqrt{[2h \cdot (6/7)/R] \cdot [1/\cos(\gamma/2) - 1]} = 1.78 \cdot \sqrt{h} \cdot [1/\cos(\gamma/2) - 1].$$

Calculations of the angle  $g = \alpha' - \alpha$  using exact formulas for several heights  $h$  (in meters) and selected fields of view  $\gamma$  (in degrees) are provided in Table 2. The absence of a percentage value in this table indicates that the entire horizon is within the field of view; in this case,  $r' = 0$ , and the chord becomes the diameter of the horizon,  $x = 2r$ .

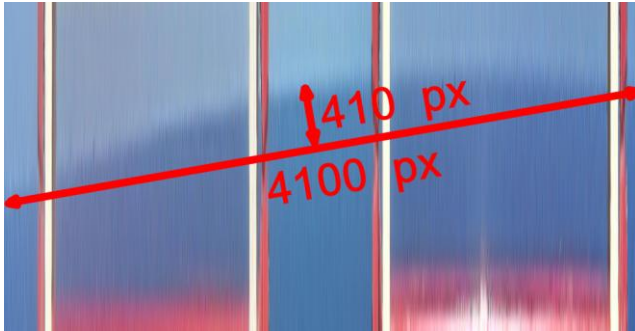
**Table 2.** Observed horizon bulge,  $g$  (expressed in arc minutes and percentage of the angle at which the chord is visible,  $\gamma$ ), above the chord as a function of the observer's height above the Earth,  $h$ , for selected fields of view,  $\gamma$ . Calculations were performed using exact formulas, taking into account mean refraction. Missing numerical data indicate situations when for very high altitudes the observer field of view envelopes entire horizon.

Height $h$	Field of View, $\gamma$									
	30°		60°		90°		120°		150°	
[m]	[']	[%]	[']	[%]	[']	[%]	[']	[%]	[']	[%]
10	0.199	0.011	0.872	0.024	2.34	0.043	5.64	0.078	16.1	0.179
100	0.629	0.035	2.759	0.077	7.39	0.137	17.8	0.248	51.1	0.567
1000	1.989	0.111	8.724	0.242	23.4	0.433	56.4	0.783	161.6	1.796
10000	6.293	0.350	27.60	0.767	73.9	1.369	178.7	2.482	514.8	5.720
25000	9.956	0.553	43.68	1.213	117.1	2.168	283.5	3.937	824.1	9.156
50000	14.09	0.783	61.86	1.718	166.0	3.074	403.1	5.598	1191.2	13.24
100000	19.97	1.109	87.72	2.437	235.9	4.368	576.4	8.006	1771.0	19.68
500000	45.34	2.519	200.6	5.573	549.7	10.18	1432.9	19.90	4172.9	—
1000000	65.32	3.629	292.0	8.111	823.6	15.31	2560.0	35.56	3708.8	—
5000000	168.2	9.344	868.7	24.13	2202.9	—	2202.9	—	2202.9	—
10000000	282.1	15.68	1514.2	—	1514.2	—	1514.2	—	1514.2	—

These calculations show that in practice, the horizon's curvature above the chord for fields of view up to 60° and an observation point altitude of up to 10 km is less than 1% of that 60° (i.e., 0.6°), making it difficult to see. Due to this small degree of horizon curvature, coupled with frequent image distortions resulting from imperfect camera optics (the "fisheye" effect), it is important that the horizon line (or rather, the chord) runs as precisely as possible through the center of the resulting image.

Flat Earth advocates claim that the lack of horizon curvature can be verified by placing a straight board or beam horizontally at eye level on the seashore. Suppose such a board is 3 meters long and placed 10 meters above sea level, and we look at it so that we can see its ends on the horizon, encompassing a horizon angle of view of 60°. This means that our eyes are at a distance of 3 m from the ends of the board, and  $3 \cdot \cos(60^\circ/2) = 2.60$  m from its center. The angle of curvature of the horizon read from Table 2 is in this case  $g = 0.872'$  (approx. 0.024% of 60°), or calculated using the given approximate formula  $1.78 \cdot \sqrt{10 \cdot [1/\cos(30^\circ) - 1]} = 0.870'$ . This further means that at a distance of the board, the horizon will protrude above it by  $2.60 \cdot \tan g = 0.0007$  m, i.e. less than 1 mm. Even if the board is perfectly level, such a small curvature is unlikely to be noticeable. A similar calculation for the field of view  $\gamma = 120^\circ$  (the distance of

the eyes from the center of the board will then be 0.87 m) gives only about twice as high bulge above the board:  $0.87 \cdot \text{tg}(5.64') = 0.0014 \text{ m}$ , i.e. 1.4 mm.



The page [Left to right curve of the horizon](#) features a photograph of the horizon, viewed through the slits of a special device, taken from a height of  $h = 184 \text{ m}$  with a field of view of  $\gamma = 63.7^\circ$  (EXIF data). The expected bulge is therefore  $1.78 \cdot \sqrt{184} \cdot [1/\cos(63.7/2) - 1] = 4.3'$ , which is  $(4.3/60)/63.7 \cdot 100 =$

0.11% of the field of view. In a fragment of the aforementioned photograph (see the accompanying image) vertically enlarged by a factor of 100, a bulge of approximately 410 px (pixels) can be seen (with an error of approximately 10%) over a chord length of approximately 4100 px, or 0.1% of the field of view. This is consistent with the expected value of 0.11%.

And one more note for flat earthers. Sometimes objects far beyond the calculated horizon are observed but it doesn't mean that the Earth isn't a sphere, as the formulas given here take into account the average astronomical refraction, and the actual refraction depends on temperature and pressure. Furthermore, one has to remember that there are phenomena such as [terrestrial refraction](#) and mirages.

K.M. Borkowski  
Toruń (Poland), 17 May 2024